

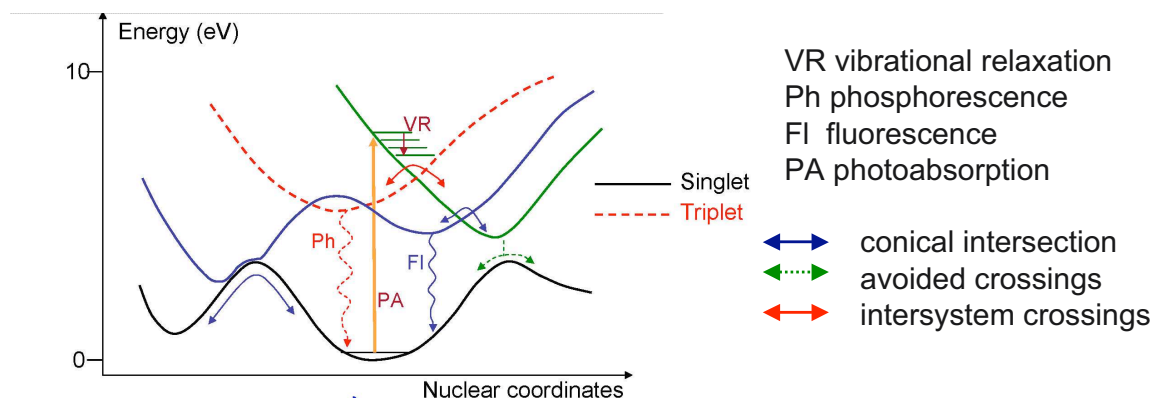
## Some Algorithms in Quantum Chemistry

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### Quantum Chemistry

- Multiple Potential Energy Surfaces (Born-Oppenheimer Approximation)
- Electron Correlation (various approximations and approaches)
- Molecular Structure and Properties (ground and excited states)
- Critical Points of high-dimensional surfaces (transition states, minima, intersections, etc.)



## COLUMBUS

**Goal:** Parallel general-purpose ab initio techniques applicable to ground and excited states for arbitrary points on the PES

**Applications:** spectroscopy, dynamics, kinetics

**History:** 1980 first sequential version, 1990 first parallel version

**Distribution:** Open source and binary (~500K lines of Fortran+C+Perl)

**Contributors:** I. Shavitt (GUGA), H. Lischka (photodynamics, solvent effects, Analytic Gradients, NAC), R. Shepard (MCSCF, MRCI, Analytic Gradients, NAC, SPAM), R. Pitzer (integrals, spin-orbit CI), P. G. Szalay (Analytic Gradients, MR-AQCC, GDIIS), Th. Mueller (parallel MRCI/MR-AQCC, interoperability, spin-orbit CI) D. R. Yarkony (NAC, Conical Intersections, crossing seams), Granucci (Surface Hopping Dynamics), and many others.

<http://www.univie.ac.at/columbus/>

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3

### Single-Point Electronic Energy and Wave Function Evaluation:

At a fixed molecular conformation and for the electronic states of interest, evaluate the expectation value

$$E(\mathbf{R}) = \frac{\langle \Psi_{el}(\mathbf{r}; \mathbf{R}) | H^{op}(\mathbf{r}; \mathbf{R}) | \Psi_{el}(\mathbf{r}; \mathbf{R}) \rangle}{\langle \Psi_{el}(\mathbf{r}; \mathbf{R}) | \Psi_{el}(\mathbf{r}; \mathbf{R}) \rangle}$$
$$H^{op} = \sum_j^{Nel} \frac{-\hbar^2}{2m_e} \nabla_j^2 + \sum_j^{Nel} \sum_a^{Natom} \frac{Z_e Z_a}{|\mathbf{r}_j - \mathbf{R}_a|} + \sum_{j < k}^{Nel} \frac{Z_e^2}{|\mathbf{r}_j - \mathbf{r}_k|}$$
$$\langle \rangle \equiv \int \int \dots \int d\mathbf{r}_1 d\mathbf{r}_2 \dots d\mathbf{r}_{Nel}$$



4

The electronic wave function is expanded in a basis:

$$|\Psi_{el}(\mathbf{r}; \mathbf{R})\rangle = \sum_m^{N_{csf}} c_m |m(\mathbf{r}; \mathbf{R})\rangle \quad \leftarrow \text{configuration state function}$$

$$|m(\mathbf{r}; \mathbf{R})\rangle = \sum_{\sigma_1 \sigma_2 \dots \sigma_{N_{el}}} c_{\sigma_1 \sigma_2 \dots \sigma_{N_{el}}} |\varphi_{m_1 \sigma_1}(\mathbf{r}_1; \mathbf{R}) \varphi_{m_2 \sigma_2}(\mathbf{r}_2; \mathbf{R}) \dots \varphi_{m_{N_{el}} \sigma_{N_{el}}}(\mathbf{r}_{N_{el}}; \mathbf{R})|$$

$$\varphi_k(\mathbf{r}; \mathbf{R}) = \sum_{\mu}^n C_{\mu k}(\mathbf{R}) \chi_{\mu}(\mathbf{r}; \mathbf{R}_{a(\mu)})$$

Slater determinant

atom-centered basis functions

orbital coefficients (optimized at each conformation)

orthonormal molecular orbitals



5

The CSF expansion results in a real symmetric eigenvalue equation:

$$\mathbf{H}(\mathbf{R})\mathbf{c}(\mathbf{R}) = \mathbf{E}(\mathbf{R})\mathbf{c}(\mathbf{R})$$

$$H(\mathbf{R})_{mn} = \langle m | H^{op} | n \rangle$$

2-particle coupling coefficient

$$= \sum_{p,q}^{Norb} h(\mathbf{R})_{pq} \langle m | E_{pq} | n \rangle + \frac{1}{2} \sum_{p,q,r,s}^{Norb} g(\mathbf{R})_{pqrs} \langle m | e_{pqrs} | n \rangle$$

2-electron hamiltonian integral

1-particle coupling coefficient

1-electron hamiltonian integral



6

## Tensor Transformations of Integral Arrays:

### Integral arrays as needed

$$h_{pq} = \sum_{\mu\nu} C_{\mu p} C_{\nu q} h_{\mu\nu}^{[ao]} = (\mathbf{C}^T \mathbf{h}^{[ao]} \mathbf{C})_{pq}$$

$$h_{\mu\nu}^{[ao]} = \int \chi_{\mu}(\mathbf{r}) \left( -\frac{1}{2} \nabla^2 - \sum_a^{N_{atom}} \frac{Z_a}{|\mathbf{r} - \mathbf{R}_a|} \right) \chi_{\nu}(\mathbf{r}) d\mathbf{r}$$

### Integral arrays as computed

$$g_{pqrs} = \sum_{\mu\nu\lambda\kappa} C_{\mu p} C_{\nu q} C_{\lambda r} C_{\kappa s} g_{\mu\nu\lambda\kappa}^{[ao]}$$

$$g_{\mu\nu\lambda\kappa}^{[ao]} = \iint \chi_{\mu}(\mathbf{r}_1) \chi_{\nu}(\mathbf{r}_1) \left( \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \right) \chi_{\lambda}(\mathbf{r}_2) \chi_{\kappa}(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

### Configuration Interaction: Real Symmetric Eigenvalue Problem

- Use the iterative Davidson Method for the lowest (or lowest few) eigenpairs
- **Direct CI:**  $\mathbf{H}$  is not explicitly constructed,  $\mathbf{w}=\mathbf{H}\mathbf{x}$  are constructed in “operator” form
- Matrix dimensions are  $10^4$  to  $10^{10}$
- All floating point calculations are 64-bit (need 8 to 11 significant figures in the energy)

## Davidson Method (simplified)

Generate an initial vector  $\mathbf{x}_1$

MAINLOOP: DO  $n=1$ , NITER

    Compute and save  $\mathbf{w}_n = \mathbf{H} \mathbf{x}_n$

    Compute the  $n^{th}$  row and column of  $\mathbf{G} = \mathbf{X}^T \mathbf{H} \mathbf{X} = \mathbf{W}^T \mathbf{X}$

    Compute the subspace Ritz pair:  $(\mathbf{G} - \rho \mathbf{1}) \mathbf{c} = 0$

    Compute the residual vector  $\mathbf{r} = \mathbf{W} \mathbf{c} - \rho \mathbf{X} \mathbf{c}$

    Check for convergence using  $|\mathbf{r}|$ ,  $\mathbf{c}$ ,  $\rho$ , etc.

    IF (converged) THEN

        EXIT MAINLOOP

    ELSE

        Generate a new expansion vector  $\mathbf{x}_{n+1}$  from  $\mathbf{r}$ ,  $\rho$ ,  $\mathbf{v} = \mathbf{X} \mathbf{c}$ , etc.

    ENDIF

ENDDO MAINLOOP



## Matrix-Vector Products

$$\mathbf{w} = \mathbf{H} \mathbf{x}$$

$$\begin{aligned} w_m &= \sum_n^{N_{\text{csf}}} H_{mn} x_n \\ &= \sum_n^{N_{\text{csf}}} \sum_{p,q}^{N_{\text{orb}}} h_{pq} \langle m | E_{pq} | n \rangle x_n + \frac{1}{2} \sum_n^{N_{\text{csf}}} \sum_{p,q,r,s}^{N_{\text{orb}}} g_{pqrs} \langle m | e_{pqrs} | n \rangle x_n \end{aligned}$$

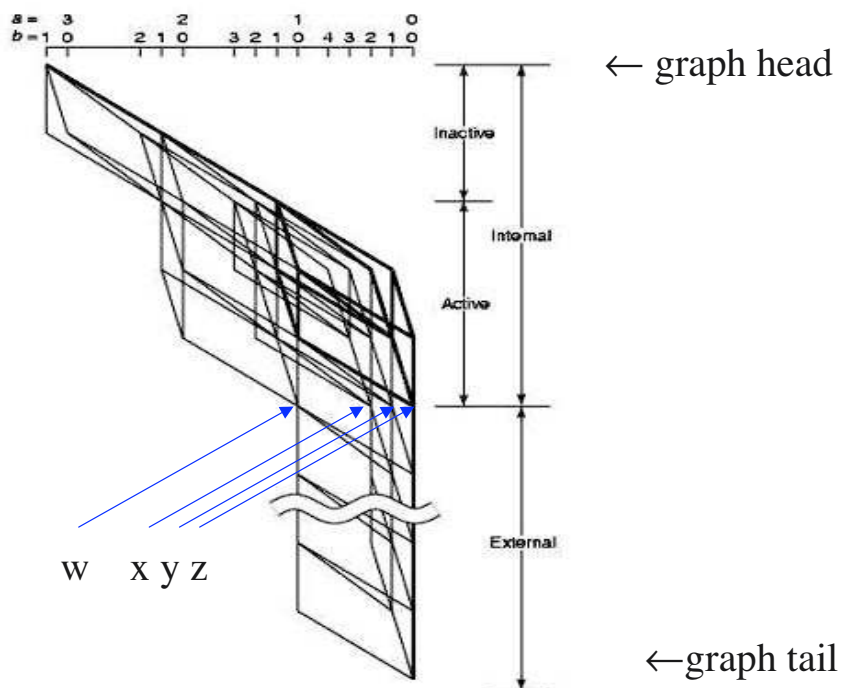
- The challenge is to bring together the different factors in order to compute  $\mathbf{w}$  efficiently
- $h_{pq}$  and  $g_{pqrs}$  are computed and stored as arrays (with index symmetry)
- $\langle m | E_{pq} | n \rangle$  and  $\langle m | e_{pqrs} | n \rangle$  are sparse and are recomputed as needed



## Coupling Coefficient Evaluation

- Graphical Unitary Group Approach (**GUGA**)
- Define a directed graph with nodes and arcs: **Shavitt Graph**
- Nodes correspond to spin-coupled states consisting of a subset of the total number of orbitals and electrons
- Arcs correspond to the (up to) four allowed spin couplings when an orbital is added to the graph
- Coupling coefficients are evaluated as products of "segment values" of interacting segments within the Shavitt Graph.

## CSF/Walk order



### H Matrix Structure

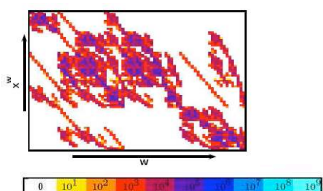
- 0:  $g_{pqrs}$
- 1:  $g_{pqra}$
- 2:  $g_{pqab}$ ,  
 $g_{pa,qb}$
- 3:  $g_{pabc}$
- 4:  $g_{abcd}$

	z	y	x	w
z	0	1	2	2
y	1	0,2	1,3	1,3
x	2	1,3	0,2,4	2
w	2	1,3	2	0,2,4

- Some integral subsets are replicated on all nodes
- Some are distributed among nodes (e.g. GA Library)

## Computational Kernels and Scaling

- Depending on the block (zz, yz, wx, etc.), the low-level arithmetic operations are DDOT, DAXPY, DGEMV, or DGEMM with the matrix dimension being  $N_{orb}$  or  $N_{ext}$ .



Relative cost for wx 2-external block in terms of groups of valid internal walks

- Total effort for MR-SDCI is  $\sim 4^{N_{active}} N_{orb}^6$

This problem can use an arbitrarily large amount of computer time as the size of the molecule increases

## Davidson Diagonalization Details

### New Expansion Vector:

$$\mathbf{x}^{[new]} = -(\mathbf{H}\mathbf{v} - \rho\mathbf{v}) = -\mathbf{r}$$

Lanczos

$$(\mathbf{H}^0 - \rho)\mathbf{x}^{[new]} = -\mathbf{r}$$

$\mathbf{H}^0 = \mathbf{D}$  → Diagonal  
Preconditioned Residual

$$(\mathbf{H}^0 - \rho)\mathbf{x}^{[new]} = -\mathbf{r} + \varepsilon\mathbf{v}$$

RQII, IIGD, GJD, etc.

$$\mathbf{x}^{[new]} = \mathbf{Q}\bar{\mathbf{v}} \quad \text{with} \quad \bar{\mathbf{H}}\bar{\mathbf{v}} = \bar{E}\bar{\mathbf{v}}$$

$$\mathbf{P} = \mathbf{X}\mathbf{X}^T; \mathbf{Q} = \mathbf{1} - \mathbf{P}; \bar{\mathbf{H}} = \mathbf{H} - \mathbf{Q}(\mathbf{H} - \mathbf{H}^0)\mathbf{Q}$$

SPAM (converging  
preconditioner)  
*J. Comp. Phys.* **172**, 472 (2001)

## ...Davidson Diagonalization Details

### Choice of $\mathbf{v}$ and $\rho$ :

- Choose the lowest unconverged eigenpair (sequential convergence)
- Cycle among unconverged eigenpairs (maintains subspace faithfulness)
- Choose the approximate eigenpair that maximizes  $\langle \mathbf{v} | \mathbf{v}^{\text{ref}} \rangle^2$  (vector following)
- Choose the approximate eigenpair that minimizes  $|\rho - E^{\text{ref}}|$  (root homing)

### Single or multiple $\mathbf{x}^{new}$ (blocked algorithms):

- This affects data reuse efficiency, task granularity, and the convergence rate

### Subspace contraction (restart strategy):

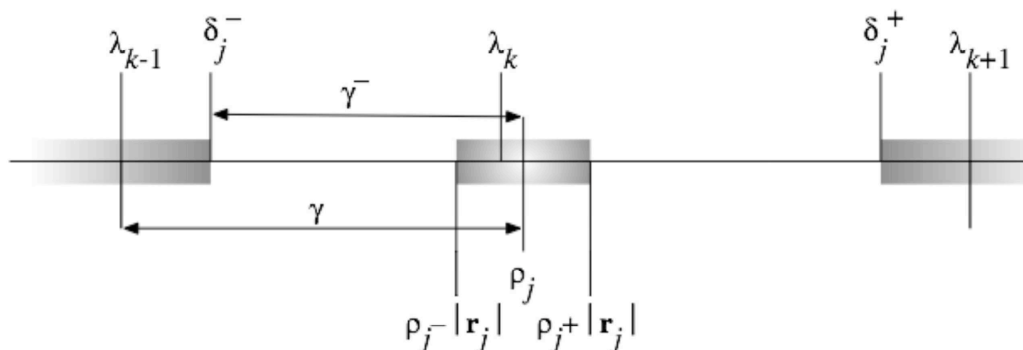
- Sometimes determined entirely by storage capacity (no choice)
- Optimal max and min subspace dimensions improves efficiency (transform away useless components) by reducing memory, I/O, and communications requirements.
- Nonoptimal dimensions hurt efficiency – sometimes difficult to predict optimal strategy



## ...Davidson Diagonalization Details

### Convergence :

- Usually do not need energies (eigenvalues) to full machine precision (8 to 11 significant figures is typical).



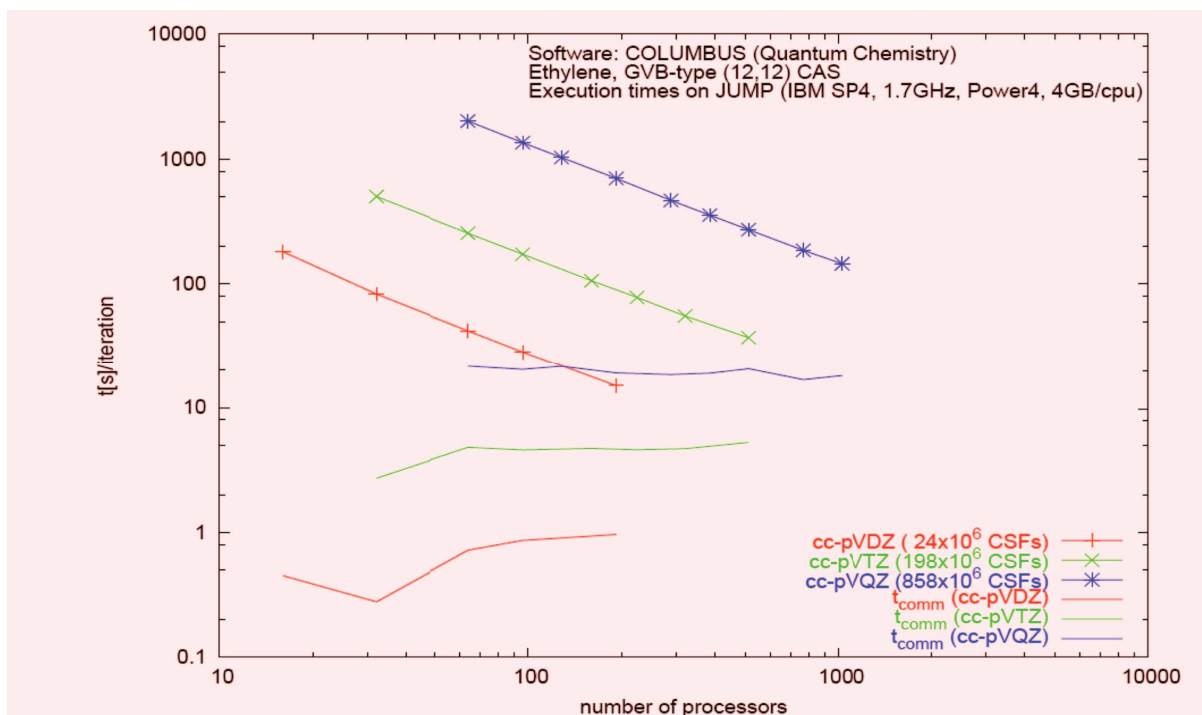
$$\frac{\mathbf{r}_j^\dagger \mathbf{M}^{-1} \mathbf{r}_j}{\gamma_j^-} \geq \frac{\mathbf{r}_j^\dagger \mathbf{M}^{-1} \mathbf{r}_j}{\gamma_j} \geq |\lambda_j - \rho_j|$$

Apply residual bound, spread bound, and gap bound recursively

*Comput. Phys. Comm.* **167**, 90 (2005)

*Comput. Phys. Comm.* **170**, 109 (2005)

## COLUMBUS Scalability



## General Convergence Acceleration: Direct Inversion in the Iterative Subspace (DIIS)

Want to solve  $\mathbf{F}(\mathbf{x}_*) = \mathbf{0}$  with  $\mathbf{F}(\mathbf{x}) \in \mathbb{R}^m$  and  $\mathbf{x} \in \mathbb{R}^q$ .

If  $\mathbf{u}^T \mathbf{c} = 1$  with  $u_k = 1$  for  $k=1:n$ , then

$$\mathbf{F}(\mathbf{x}) + \mathcal{O}(\|\mathbf{x} - \mathbf{x}_0\|^2) = \sum_{k=1}^n \mathbf{F}(\mathbf{x}_k) c_k = \sum_{k=1}^n \mathbf{e}_k c_k = \mathbf{E} \mathbf{c}$$

Least squares interpolation is given by

$$\min_{\mathbf{c}} \|\mathbf{E} \mathbf{c}\| \quad \text{with } \mathbf{u}^T \mathbf{c} = 1$$

$$\mathbf{x}^{\text{DIIS}} = \mathbf{X} \mathbf{c}$$

## ...DIIS

$$\begin{pmatrix} \mathbf{B} & -\mathbf{u} \\ -\mathbf{u}^T & 0 \end{pmatrix} \begin{pmatrix} \mathbf{c} \\ \lambda \end{pmatrix} = \begin{pmatrix} \mathbf{0} \\ -1 \end{pmatrix}$$

Normal equation

$$\min_{\mathbf{c}} \|\mathbf{E} \mathbf{c} - \mathbf{b}\| \quad \text{with } \mathbf{D} \mathbf{c} = \mathbf{d}$$

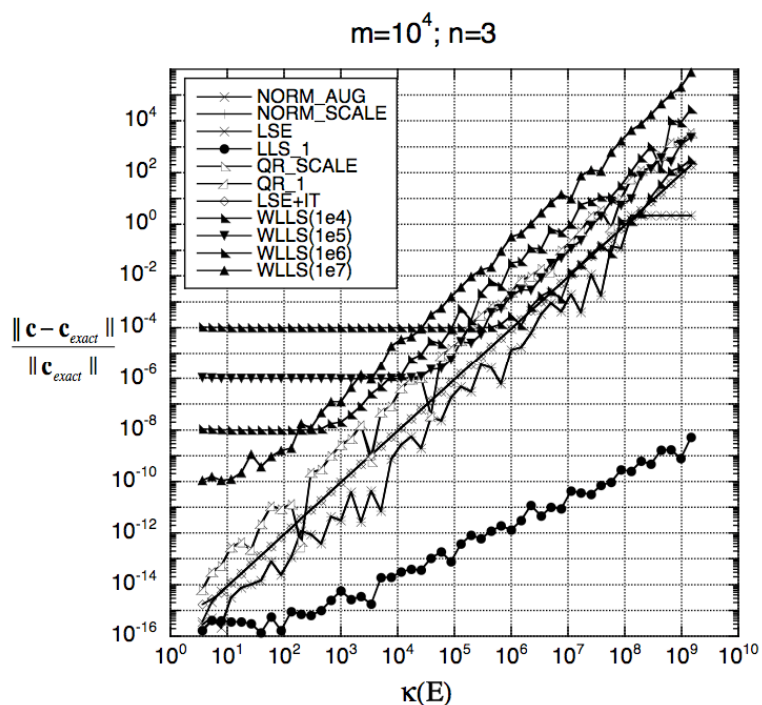
LSE

$$\min_{\tilde{\mathbf{c}}} \|\tilde{\mathbf{E}} \tilde{\mathbf{c}} + \mathbf{e}_n\|$$

Substitution and elimination

$$\min_{\mathbf{c}} \left\| \begin{pmatrix} \mathbf{E} \\ w \mathbf{u}^T \end{pmatrix} \mathbf{c} - \begin{pmatrix} \mathbf{0} \\ w \end{pmatrix} \right\|$$

Weighted least squares



## Summary

### Configuration Interaction

- Seek efficient ways to distribute  $\mathbf{h}$  and  $\mathbf{g}$  arrays on parallel computers
- Seek efficient and flexible iterative subspace diagonalization methods (need 8 to 11 significant figures in the eigenvalues)
- Consider contractions with  $\mathbf{h}^{[AO]}$  and  $\mathbf{g}^{[AO]}$  – more arithmetic operations per task, but no storage or communications requirements
- Tasks are not uniform in size – this presents load balancing challenges

### Tensor Transformation

- Need efficient and general tensor transformation and tensor contraction libraries
- Must include index symmetry to reduce storage and arithmetic operation counts

### Least Squares Interpolation (DIIS)

- Need robust software with  $\epsilon\kappa(E)$  error rather than  $\epsilon\kappa(E)^2$  error